

278 short-distance dynamics) is computed using the colour-singlet model, where the wave func-

279 ~~tion at the origin $|R(0)|^2 = 0.48 \text{ GeV}^3$ is determined from the leptonic decay width of the J/ψ :~~

irrelevant.

280 ~~$\Gamma(J/\psi \rightarrow e^+ e^-) = \frac{16\alpha^2}{9m_\psi^2} |R(0)|^2 = 4.72 \text{ keV}$.~~

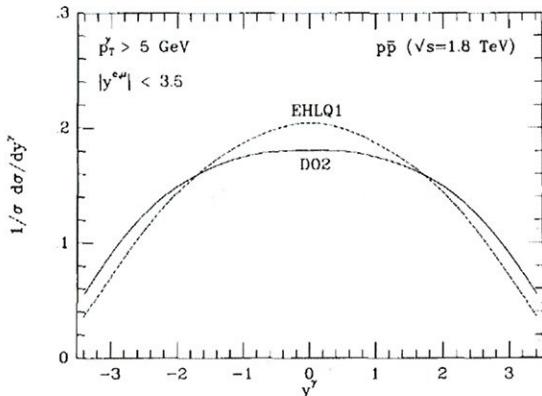


FIG. 4. Normalised differential cross section in y for $J/\psi + \gamma$ (taken from [4]).

281 $\tilde{p}p$ collisions \rightarrow unlabelled, or boldface

282 The cuts applied to the data are the following: $p_T^\gamma = p_T^\psi > 5 \text{ GeV}$ for the photon and J/ψ trans-

In order to match the experimental constraints,

283 verse momentum and $|y^{\gamma,e,\mu}| < 3.5$ for the rapidities of the photon and the leptons. In Fig. 4 the

284 rapidity distribution and the p_T spectrum of the photon are displayed, for two different parametri-

285 sations of the gluon content of the photon, namely the EHLQ1 [?] and D02 [?], which make

Despite the different cut

286 different assumptions about the behaviour of the $f_{g/p}$ structure functions. Despite their different

Nowadays, the knowledge of gluon PDF is of course better and one can systematically study the uncertainty

287 behaviour at large and small x , the two parametrisations lead to similar results for the differential

of the PDF on given observable. It would therefore be very interesting to update this analysis with

288 cross section, as shown. Therefore, although the events produced at $\tilde{p}p$ collisions might not allow

modern theoretical tools.

289 to extract useful information about the normalisation of the gluon densities $f_{g/p}$, they might allow

290 a detailed measurement of their shape at small x , down to $\sim 10^{-4}$.

291 $e p$ collisions

292 In this case the following cuts have been applied: $p_T^\gamma = p_T^\psi > 1.5 \text{ GeV}$ and $-3.5 < y^{\gamma,e,\mu} < 3$
 293 (negative rapidities correspond to the proton beam direction). An observation of a sizable signal,
 294 in these measurements, would be ~~an unambiguous proof for the existence of the gluon content~~
 295 inside the photon. The gluon density functions of the photon and the electron are $f_{g/\gamma}$ and $f_{g/e}$)?
 296 respectively. For the calculation of the latter, the gluon content of the photon $f_{g/\gamma}$ is convoluted
 297 with the photon content of the electron $f_{\gamma/e}$. However, there does not exist a momentum sum rule
~~the gluon density in a photon,~~
 298 for $f_{g/\gamma}$, to allow to extract information on its shape and renormalisation. Fig? shows the nor-
~~the authors studied~~
 299 malised rapidity distribution of the photon, using three different parametrisations of the photon
 300 structure functions: DG, LAC1 and LAC3. The first one, ~~is~~ based on the assumption that the gluon
 301 densities inside photons come from gluon radiation off quarks, while in the analysis undertaken
~~and the others assuming~~
 302 by the other two, the existence of an "intrinsic" gluon content of the photon is taken into account.
~~which they obtained shows the discriminatory power of J/ψ + γ on f_{g/\gamma}.~~
 303 The large differences between the three curves in fig? account for the lack of data constraining
~~As for the pp analysis, it would be worth updating such a study with much modern tools~~
 304 $f_{g/\gamma}$, pointing to the essence of more $J/\psi + \gamma$ events at HERA. It should be stated, though, that
~~in view of the Electron - Ion Colliders and LHeC projects.~~
 305 there is no guarantee that data will be described by any of these parametrisations.

CONCLUSION : Recall the motivations and give your feelings

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