





AFTER@LHC: A fixed-target programme at the LHC for heavy-ion, hadron, spin and astroparticle physics

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AFTER@LHC Study group: http://after.in2p3.fr/after/index.php/Current_author_list



Part I

Why a new fixed-target experiment for High-Energy Physics now?

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- which are essential assets to study
 - rare proton fluctuations at large *x*
 - vector boson production near threshold and other rare processes
 - nuclear dependence in heavy-ion collisions
 - observables involving gluons and the target proton spin

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 - Very large PDF uncertainties for $x \gtrsim 0.5$.

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- Proton charm content important to high-energy neutrino & cosmic-rays physics
- EMC effect is an open problem; studying a possible gluon EMC effect is essential
- Relevance of nuclear PDF to understand the initial state of heavy-ion collisions
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- Explore the longitudinal expansion of QGP formation with new hard probes
- Test the factorisation of cold nuclear effects from p + A to A + B collisions
- Test the formation of azimuthal asymmetries: hydrodynamics vs. initial-state radiation

Part II

A fixed-target experiment using the LHC beam(s): AFTER@LHC

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- Let us simply avoid the forward region! How?



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backward physics = $large-x_2$ physics

Part III

Colliding the LHC beams on fixed targets: 2 options

★ The LHC beam may be extracted using "Strong crystalline field"

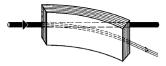
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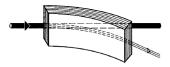
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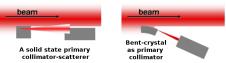
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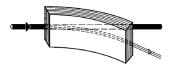
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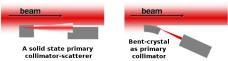
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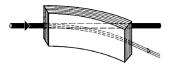
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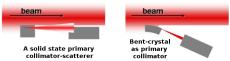
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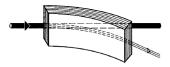
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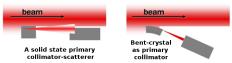
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- ★ CRYSBEAM: ERC funded project to extract the LHC beams
 - with a bent crystal (G. Cavoto Rome)

Luminosities with extracted-proton beams

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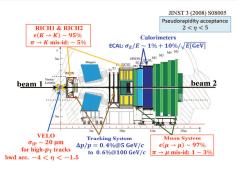
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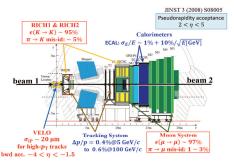
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3 orders of magnitude larger than RHIC (200 GeV)





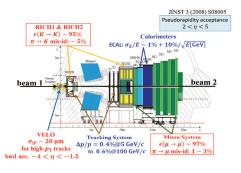
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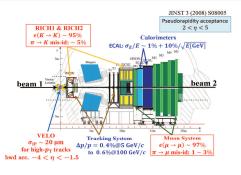
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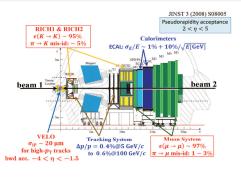
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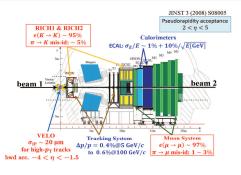
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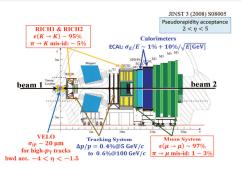
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- Target density: $\frac{\rho}{P} = c = \frac{A}{22400} \text{bar}^{-1} g \, cm^{-3} \Rightarrow \mathcal{L} = \Phi_{beam} \times (\frac{\mathcal{N}_A}{22400} \times P \times \ell)$

[1 mole of a perfect gas occupies 22 400 \mbox{cm}^3 at 273 K and 1 bar]

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- Target density: $\frac{\rho}{P} = c = \frac{A}{22400} \text{bar}^{-1} g \, cm^{-3} \Rightarrow \mathcal{L} = \Phi_{beam} \times \left(\frac{\mathcal{N}_A}{22400} \times P \times \ell\right)$

[1 mole of a perfect gas occupies 22 400 cm³ at 273 K and 1 bar]

• For $P = 10^{-9}$ bar [7× that of SMOG in 2015, the 'vacuum' is 10^{-12} bar], $\mathcal{L}_{DX(PbX)} = 10(10^{-3})\mu b^{-1} s^{-1}$

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which could be polarised

C. Barschel, P. Lenisa, A. Nass, and E. Steffens, Adv.Hi.En.Phys. (2015) 463141; See E. Steffens's talk at PSTP 2015

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- \Rightarrow For PbA, limitations would come first from the beam lifetime, pile-up and exp. DAQ
- A specific gas target is a competitive alternative to the beam extraction

[1/2 Ampère !]

Luminosities with a polarised internal-gas-target option

Advances in High Energy Physics Volume 2015, Article ID 463141, 6 pages http://dx.doi.org/10.1155/2015/463141

A Gas Target Internal to the LHC for the Study of pp Single-Spin Asymmetries and Heavy Ion Collisions

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¹LHCb Collaboration, CERN, 1211 Geneva 23, Switzerland

²University of Ferrara and INFN, 44100 Ferrara, Italy ³Institut für Kernphysik, FZJ, 52425 Jülich, Germany

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$$\int dt \mathcal{L} = 10^{33} \text{cm}^{-2} s^{-1} \stackrel{\Delta t = 10^7 \text{ s}}{=} 10 \text{ fb}^{-1}!$$

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The authors claim that the COMPASS type frozen spin target machinery takes too much space in the LHC tunnel. Instead, a UVa-type NH₃ DNP target* with smaller target set-up may be considered for this comparison with parameters§:

$$n_t = 1.5 \ 10^{23} / cm^2$$
, $P_p = 0.85$, dilution $f = 0.17$.

This results in a FoM = $n_t P^2 f^2 = 3.1 \cdot 10^{21}$ /cm². As the beam intensity i_p also enters the measurement quality, we define

$$FoM* = i_p \cdot FoM = P^2 \cdot f^2 \cdot i_p \cdot n_t = P^2 \cdot f^2 \cdot \mathcal{L}$$

E. Steffens's talk at PSTP 2015

"The authors"= we "claim" = in fact, we believe so, but may be mislead

Results:

UVa-target and bent-crystal extr. beam

OMPASS-target " " "

'HERMES' target and full LHC beam (T = 300/100 K, P = 0.85, α = 0.95)

§) note that $n_t = \text{density of target nucleons}$; then $f \cdot n_t$ is the number of polarizable nucleons

$$FoM* = 1.57 \cdot 10^{30}/cm^2 s$$

$$foM* = 1.87 \cdot 10^{32}/cm^2 s$$

 $FoM* = 0.60/1.04 \cdot 10^{33}/cm^2 s$



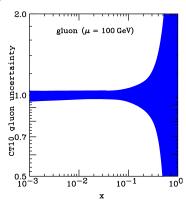
13 / 37

Part IV

AFTER@LHC: [a selection of] the physics case

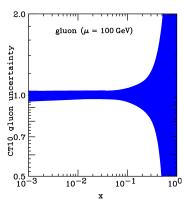
• Gluon distribution at mid, high and ultra-high *x* in the proton

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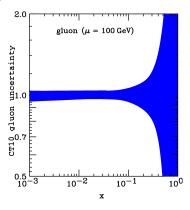
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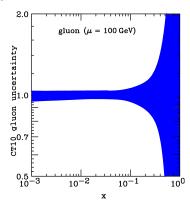
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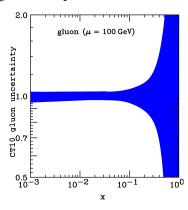
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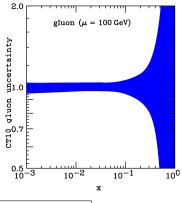
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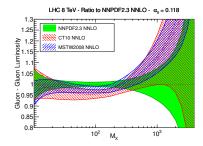
Multiple probes needed to check factorisation



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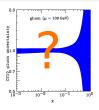
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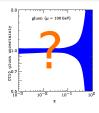
Large-*x* gluons: important to characterise some possible BSM findings at the LHC

pd physics: gluons in the neutron and the deuteron



Gluon PDF for the neutron unknown

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Gluon PDF for the neutron unknown possible experimental probes

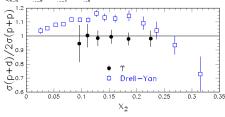
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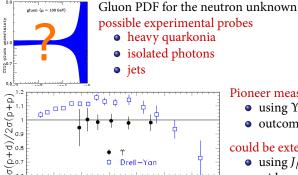
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- Pioneer measurement by E866 using $\Upsilon \rightarrow Q^2 \simeq 100 \text{ GeV}^2$
 - outcome: $g_n(x) \simeq g_p(x)$

pd physics: gluons in the neutron and the deuteron

0.3



- isolated photons

Drell-Yan

X2

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could be extended with AFTER

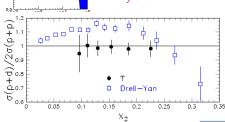
- using J/ψ , ..., C = +1 onia, ...
- wider x range & lower Q²

pd physics: gluons in the neutron and the deuteron



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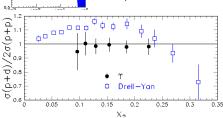
target	yearly lumi	$\mathcal{B} rac{dN_{J/\psi}}{dy}$	$\mathcal{B} rac{dN_{\Upsilon}}{dy}$
1m Liq. H ₂	20 fb ⁻¹	4.0×10^{8}	9.0×10^{5}
1m Liq. D ₂	24 fb^{-1}	9.6×10^{8}	1.9×10^{6}

pd physics: gluons in the neutron and the deuteron



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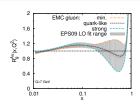
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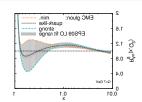
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If $g_n(x) - g_p(x)$ is too small, this measurement would anyhow be sensitive to the EMC and Fermi-motion effects in the deuteron

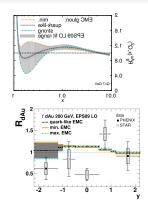
- Large-*x* gluon nPDF: unknown
- Gluon EMC effect: unknown



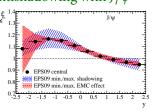
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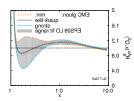


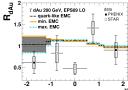
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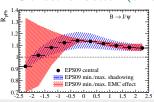
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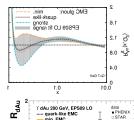


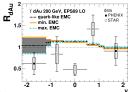




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- Quest for the gluon EMC effect for bottom(onium)







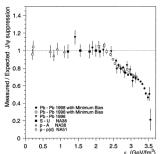
L. Massacrieret al., Adv. Hi. En. Phys. (2015) 986348

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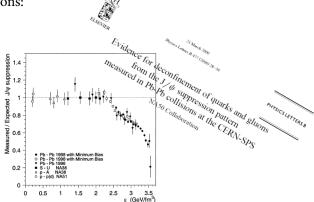
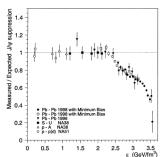


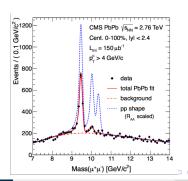
Fig. 7. Measured J/ψ production yields, normalised to the yields expected assuming that the only source of suppression is the

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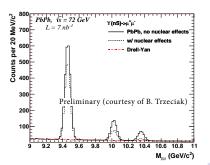
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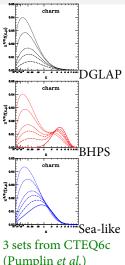
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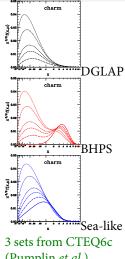
Simulations done using generic LHCb performances (incl. acceptance and efficiencies)

• Heavy-quark distributions (at high x_B)

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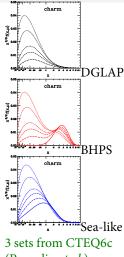


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 - c(x) & b(x) & the 5-flavour scheme at large x for BSM studies F.Maltoni,..., JHEP 1207 (2012) 022



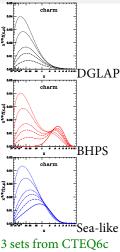
(Pumplin et al.)

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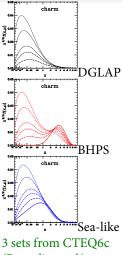
(Pumplin et al.)

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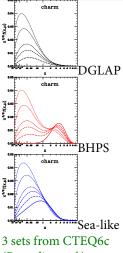
3 sets from CTEQ66 (Pumplin *et al.*)

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 - good coverage in the target-rapidity region



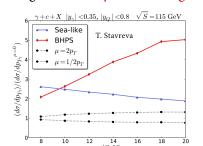
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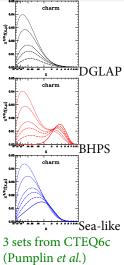
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 Uncertainties in atmospheric neutrino flux (background of cosmic neutrinos) dominated by those on charmed meson decays

IceCube collab. PRL 111 (2013) 021103; Science 342 (2013) 1242856

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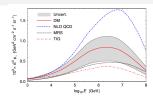


FIG. 6 (color online). Prompt muon neutrino fluxes oblanied in perturbative QCD. The shaded area represents the theoretical uncertainty in the prompt neutrino flux evaluated in this paper, and the solid line in the band is our standard result. The dished curve is the NLO perturbative QCD calculation of Ref. [14] (PSB, modified here is include fragmentation; the dotted curve is the suternation model; result of Ref. [19] (MSS); and the dashcritical perturbative QCD calculation of Ref. [15] (TG).

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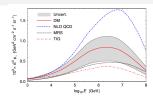


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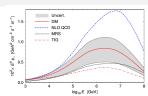


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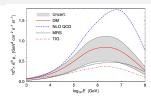


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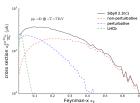


Figure 8. Weighted spectrum for D-mesons in SIBYLL at $\sqrt{s} = 7$ TeV. The contributions from the perturbative and non-perturbative model components are shown by the blue and red lines, respectively. Note the negligible contribution to the energy spectrum from the phase space covered by the LHCb experiment (2.5 < w < 4.5 green line).

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$$\begin{array}{l} x_F^{collider} = \frac{2m_T}{2E_{beam}} \; \text{sinh} \left(y^{lab.}\right) \; ; \; x_F^{FT} = \frac{2m_T}{\sqrt{2m_N E_{beam}}} \; \text{sinh} \left(y^{lab.} - 4.8\right) \\ x_F^{FT} \left(P_D^D = 0, y^{lab.} = 2\right) \simeq -0.2 \; ; \; x_F^{FT} \left(P_D^D = 4 \text{GeV}, y^{lab.} = 2\right) \simeq -0.6 \\ \end{array}$$



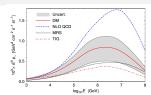


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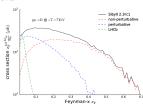


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 Similar conclusion for the ALICE muon spectrometer

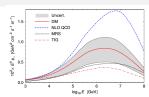


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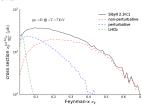


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The quest for the orbital angular momentum of the quarks and gluons

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• Quark/Gluon Sivers function: distortion in the distribution of an unpolarised partons with momentum fraction x and transverse momentum k_{\perp} due to the proton transverse polarisation: $f_{1T}^{1}(x, \vec{k}_{\perp}^2)$

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- Several experiments wish to measure $A_N^{Drell-Yan}$ to extract $f_{1T}^{\perp q}(x, \vec{k}_{\perp}^2)$
 - COMPASS: valence quarks using a pion beam (160 GeV)

on a polarised proton target

- E1027: valence quarks using a polarised proton beam (120 GeV)
 - on an unpolarised proton target
- E1039: sea quarks using an unpolarised proton beam (120 GeV)
 on a polarised proton target

Expected asymmetries

The target-rapidity region (negative x_F) corresponds to high x^{\uparrow} where the k_T -spin correlation is the largest

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The target-rapidity region (negative x_F) corresponds to high x^{\uparrow} where the k_T -spin correlation is the largest How large?

Azimuthal asymmetries in lepton-pair production at a fixed-target experiment using the LHC beams (AFTER)

Tianbo Liu¹, Bo-Qiang Ma^{1,2,a}

Transverse Single-Spin Asymmetries in Proton-Proton Collisions at the AFTER@LHC Experiment in a TMD Factorisation Scheme

M. Anselmino, 1,2 U. D'Alesio, 3,4 and S. Melis 1

¹School of Physics and State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing 100871, China
²Center for High Energy Physics, Peking University, Beijing 100871, China

¹Dipartimento di Fisica, Università di Torino, Via P. Giuria 1, 10125 Torino, Italy

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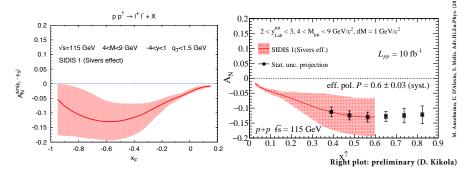
³Dipartimento di Fisica, Università di Cagliari, Cittadella Universitaria, 09042 Monserrato, Italy

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Expected asymmetries

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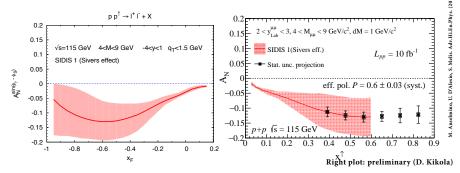
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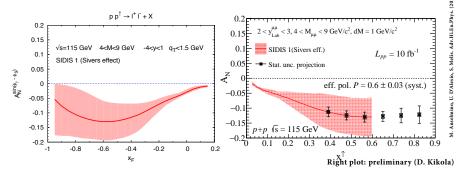


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- With 10 fb⁻¹, one can indeed expect up to 10^6 DY events in 4 < M < 9 GeV
- W and Z should be reachable with 10 fb⁻¹: $x^{\uparrow} \simeq 0.7 \div 0.8$

Part V

First simulation results

B. Trzeciak, L. Massacrier et al., Adv.Hi.En.Phys. (2015) 986348

B. Trzeciak, L. Massacrier et al., Adv. Hi. En. Phys. (2015) 986348

• LHCb has successfully carried out pPb and Pbp analyses at 5 TeV

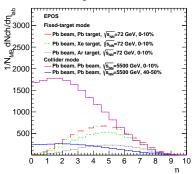


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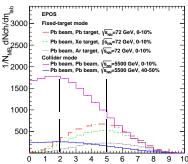
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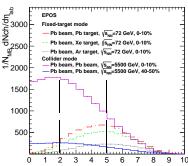
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- Simulation backed-up with a comparison of the number-of-track distribution between simulations at the detector level and data

Fast simulation using LHCb reconstruction parameters

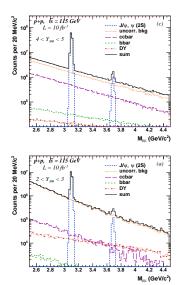
Projection for a LHCb-like detector

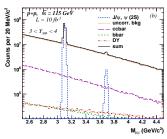
L. Massacrier, B. Trzeciak, et al., Adv.Hi.En.Phys. (2015) 986348

- Simulations with Pythia 8.185
- the LHCb detector is NOT simulated but LHCb reconstruction parameters are introduced in the fast simulation (resolution, analysis cuts, efficiencies,...)
- Requirements:
 - Momentum resolution : $\Delta p/p = 0.5\%$
 - Muon identification efficiency: 98%
- Cuts at the single muon level
 - $2 < \eta_{\mu} < 5$
 - $p_{T\mu} > 0.7 \text{ GeV}$
- Muon misidentification:
 - If π and K decay before the calorimeters (12m), they are rejected by the tracking
 - otherwise a misidentification probability is applied following: F. Achilli et al, arXiv:1306.0249

Charmonium background & its rapidity dependence

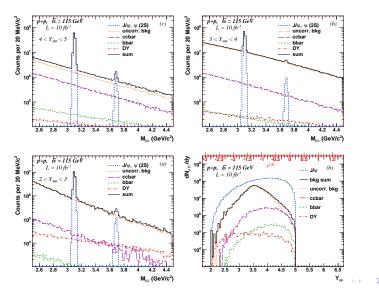
B. Trzeciak, L. Massacrier et al., 1504.05145 [hep-ex], Adv. Hi. En. Phys. (2015) 986348





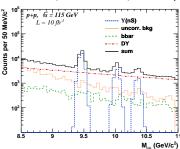
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Bottomonium background & signal reach

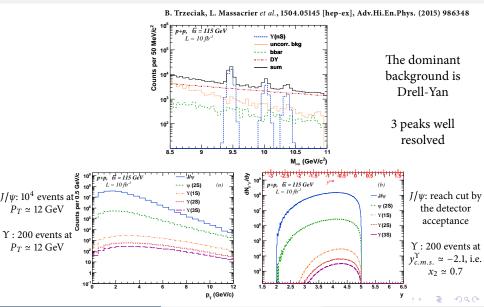
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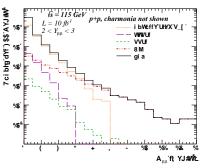
The dominant background is Drell-Yan

3 peaks well resolved

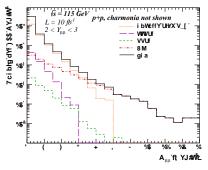
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 $\bullet \ \, \text{At backward rapidities, quark-induced processes are favoured} \Rightarrow \text{Bkgd get smaller} \\$

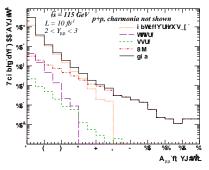


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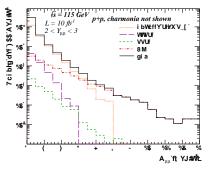
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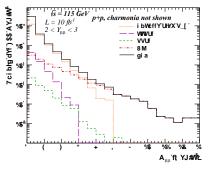
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Part VI

Further readings

Heavy-Ion Physics

- Gluon shadowing effects on J/ψ and Y production in p+Pb collisions at $\sqrt{s_{NN}}$ = 115 GeV and Pb+p collisions at $\sqrt{s_{NN}}$ = 72 GeV at AFTER@LHC by R. Vogt. Adv.Hi.En.Phys. (2015) 492302.
- Prospects for open heavy flavor measurements in heavy-ion and p+A collisions in a fixed-target experiment at the LHC by D. Kikola. Adv.Hi.En.Phys. (2015) 783134
- Quarkonium suppression from coherent energy loss in fixed-target experiments using LHC beams by F. Arleo, S.Peigné. [arXiv:1504.07428 [hep-ph]]. Adv.Hi.En.Phys. (2015) 961951
- Anti-shadowing Effect on Charmonium Production at a Fixed-target Experiment Using LHC Beams by K. Zhou, Z. Chen, P. Zhuang. Adv. High Energy Phys. 2015 (2015) 439689
- Lepton-pair production in ultraperipheral collisions at AFTER@LHC
 By J.P. Lansberg, L. Szymanowski, J. Wagner. JHEP 1509 (2015) 087
- Quarkonium Physics at a Fixed-Target Experiment using the LHC Beams. By J.P. Lansberg, S.J. Brodsky, F. Fleuret, C. Hadjidakis. [arXiv:1204.5793 [hep-ph]]. Few Body Syst. 53 (2012) 11.

Spin physics

- Transverse single-spin asymmetries in proton-proton collisions at the AFTER@LHC experiment by K. Kanazawa, Y. Koike, A. Metz, and D. Pitonyak. [arXiv:1502.04021 [hep-ph]. Adv.Hi.En.Phys. (2015) 257934.
- Transverse single-spin asymmetries in proton-proton collisions at the AFTER@LHC experiment in a
 TMD factorisation scheme by M. Anselmino, U. D'Alesio, and S. Melis. [arXiv:1504.03791 [hep-ph]].
 Adv.Hi.En.Phys. (2015) 475040.
- The gluon Sivers distribution: status and future prospects by D. Boer, C. Lorcé, C. Pisano, and J. Zhou. [arXiv:1504.04332 [hep-ph]]. Adv.Hi.En.Phys. (2015) 371396
- Azimuthal asymmetries in lepton-pair production at a fixed-target experiment using the LHC beams (AFTER) By T. Liu, B.Q. Ma. Eur.Phys.J. C72 (2012) 2037.
- Polarized gluon studies with charmonium and bottomonium at LHCb and AFTER By D. Boer, C. Pisano. Phys.Rev. D86 (2012) 094007.

Hadron structure

- Double-quarkonium production at a fixed-target experiment at the LHC (AFTER@LHC).
 by J.P. Lansberg, H.S. Shao. [arXiv:1504.06531 [hep-ph]]. To appear in Nucl. Phys. B
- Next-To-Leading Order Differential Cross-Sections for Jpsi, psi(2S) and Upsilon Production in Proton-Proton Collisions at a Fixed-Target Experiment using the LHC Beams (AFTER@LHC) by Y. Feng, and J.X. Wang. Adv.Hi.En.Phys. (2015) 726393.
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- A review of the intrinsic heavy quark content of the nucleon
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Feasibility study and technical ideas

- Feasibility studies for quarkonium production at a fixed-target experiment using the LHC proton and lead beams (AFTER@LHC) by L. Massacrier, B. Trzeciak, F. Fleuret, C. Hadjidakis, D. Kikola, J.P.Lansberg, and H.S. Shao arXiv:1504.05145 [hep-ex]. Adv.Hi.En.Phys. (2015) 986348
- A Gas Target Internal to the LHC for the Study of pp Single-Spin Asymmetries and Heavy Ion Collisions by C. Barschel, P. Lenisa, A. Nass, and E. Steffens. Adv.Hi.En.Phys. (2015) 463141
- Quarkonium production and proposal of the new experiments on fixed target at LHC by N.S. Topilskaya, and A.B. Kurepin. Adv.Hi.En.Phys. (2015) 760840

Generalities

Physics Opportunities of a Fixed-Target Experiment using the LHC Beams
 By S.J. Brodsky, F. Fleuret, C. Hadjidakis, J.P. Lansberg. [arXiv:1202.6585 [hep-ph]]. Phys.Rept. 522 (2013) 239.



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Journal Menu

- About this Journal

Abstracting and Indexing - Advance Access - Aims and Scope · Annual Issues

- Article Processing Charges Articles in Press Author Guidelines
- Bibliographic Information Citations to this Journal Contact Information Editorial Board
- Editorial Workflow - Free eTOC Alerts Publication Ethics
- Reviewers Acknowledgment Submit a Manuscript
- Subscription Information · Table of Contents

II Onen Sperial Issues @ Published Special Issues Special Issue Guidelines

Physics at a Fixed-Target Experiment Using the LHC Beams

Guest Editors: Jean-Philippe Lansberg, Gianluca Cavoto, Cynthia Hadiidakis, Jibo He, Cédric Lorcé, and Barbara Trzeciak

- Physics at a Fixed-Target Experiment Using the LHC Beams, Jean-Philippe Lansberg, Gianluca Cavoto, Cynthia Hadjidakis, Jibo He, Cédric Lorcé, and Barbara Trzeciak
 - Volume 2015 (2015). Article ID 319654. 2 pages ▶ Next-to-Leading Order Differential Cross Sections for 1/w, w(2S), and Y
- Production in Proton-Proton Collisions at a Fixed-Target Experiment Using the LHC Beams, Yu Feng and Jian-Xiong Wang Volume 2015 (2015), Article ID 726393, 7 pages
- ▶ The Gluon Sivers Distribution: Status and Future Prospects, Daniël Boer, Cédric Lorcé, Cristian Pisano, and Jian Zhou Volume 2015 (2015), Article ID 371396, 10 pages
- ▶ Studies of Backward Particle Production with a Fixed-Target Experiment Using the LHC Beams, Federico Alberto Ceccopieri Volume 2015 (2015), Article ID 652062, 9 pages
- Bremsstrahlung from Relativistic Heavy Ions in a Fixed Target Experiment at the LHC, Rune E. Mikkelsen, Allan H. Sørensen, and Ulrik I. Uggerhøj Volume 2015 (2015), Article ID 625473, 4 pages
- Antishadowing Effect on Charmonium Production at a Fixed-Target Experiment Using LHC Beams, Kal Zhou, Zhengyu Chen, and Pengfei Zhuang Volume 2015 (2015), Article ID 439689, 8 pages
- ▶ Quarkonium Production and Proposal of the New Experiments on Fixed Target at the LHC, A. B. Kurepin and N. S. Topilskava Volume 2015 (2015), Article ID 760840, 13 pages
- François Arleo and Stéphane Peigné
- ▶ Transverse Single-Spin Asymmetries in Proton-Proton Collisions at the AFTER@LHC Experiment in a TMD Factorisation Scheme, M. Anselmino, U. D'Alesio, and S. Melis Volume 2015 (2015), Article ID 475040, 12 pages
- ▶ Transverse Single-Spin Asymmetries in Proton-Proton Collisions at the AFTER@LHC Experiment, K. Kanazawa, Y. Koike, A. Metz, and D. Pitonyak Volume 2015 (2015), Article ID 257934, 9 pages
- Feasibility Studies for Quarkonium Production at a Fixed-Target Experiment Using the LHC Proton and Lead Beams (AFTER@LHC), L. Massacrier, B. Trzeciak, F. Fleuret, C. Hadjidakis, D. Kikola, J. P. Lansberg, and H.-S. Shao
- Volume 2015 (2015), Article ID 986348, 15 pages

Volume 2015 (2015). Article ID 961951, 6 pages

- ▶ Gluon Shadowing Effects on J/ψ and Y Production in p + Pb Collisions at √S_{NN} = 115 GeV and Pb + p Collisions at \(\sigma_{\text{SUN}} = 72\) GeV at AFTER@LHC, R. Vogt Volume 2015 (2015), Article ID 492302, 10 pages
- ▶ Prospects for Open Heavy Flavor Measurements in Heavy Ion and p + A Collisions in a Fixed-Target Experiment at the LHC, Daniel Kikola Volume 2015 (2015), Article ID 783134, 8 pages
- A Gas Target Internal to the LHC for the Study of pp Single-Spin Asymmetries and Heavy Ion Collisions. Colin Barschel, Paolo Lenisa, Alexander Nass, and Erhard Steffens Volume 2015 (2015). Article ID 463141. 6 pages
- A Review of the Intrinsic Heavy Quark Content of the Nucleon, S. J. Brodsky, A. Kusina, F. Lyonnet, I. Schienbein, H. Spiesberger, and R. Vogt Volume 2015 (2015), Article ID 231547, 12 pages

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Physics Reports 522 (2013) 239-255



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Physics opportunities of a fixed-target experiment using LHC beams

S.J. Brodsky a, F. Fleuret b, C. Hadjidakis c, J.P. Lansberg c,*

5.3. Color filtering, energy loss, Sudakov suppression and hadron break-up in the nucleus

a SLAC National Accelerator Laboratory, Stanford University, Menlo Park, CA 94025, USA

b Laboratoire Leprince Ringuet, Ecole polytechnique, CNRS/IN2P3, 91128 Palaiseau, France

CIPNO, Université Paris-Sud, CNRS/IN2P3, 91406 Orsav, France

Co

ntents				
ı.	Introduction	Decon	finement in heavy-ion collisions	
2.	Key numbers and features	6.1.	Quarkonium studies	
8.	Nucleon partonic structure	6.2.	Jet quenching	
	3.1. Drell-Yan	6.3.	Direct photon	
	3.2. Gluons in the proton at large x	6.4.	Deconfinement and the target rest frame	
	3.2.1. Quarkonia	6.5.	Nuclear-matter baseline	
	3.2.2. Jets	W and	1Z boson production in pp, pd and pA collisions	
		7.1.	First measurements in pA	
	3.2.3. Direct/isolated photons	7.2.	W /Z production in pp and pd	
	3.3. Gluons in the deuteron and in the neutron 8	Exclus	ive, semi-exclusive and backward reactions	
	3.4. Charm and bottom in the proton	8.1.	Ultra-peripheral collisions	
	3.4.1. Open-charm production	8.2.	Hard diffractive reactions	
	3.4.2. $J/\psi + D$ meson production	8.3.	Heavy-hadron (diffractive) production at $x_F \rightarrow -1$	
	3.4.3. Heavy-quark plus photon production	8.4.	Very backward physics	
	Spin physics	8.5.	Direct hadron production	
	4.1. Transverse SSA and DY	Furthe	er potentialities of a high-energy fixed-target set-up	
	4.2. Quarkonium and heavy-quark transverse SSA	9.1.	D and B physics	
	4.3. Transverse SSA and photon	9.2.	Secondary beams	
	4.4. Spin asymmetries with a final state polarization	9.3.	Forward studies in relation with cosmic shower	
i	Nuclear matter		usions	
			Acknowledgments	
5.2.1. Isolated photons and photon–jet correlations				
	5.2.2. Precision quarkonium and heavy-flavour studies			

Part VII

Conclusion and outlooks

Conclusion

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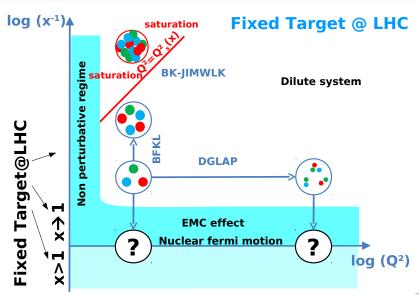
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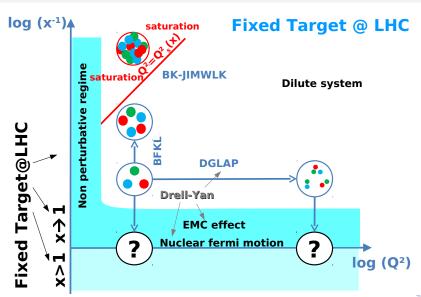
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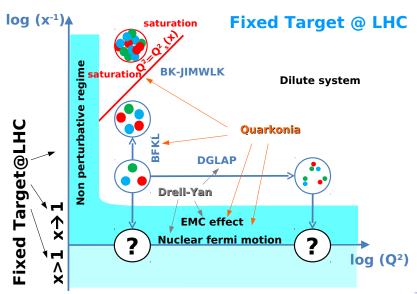


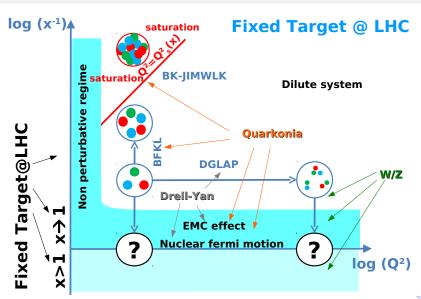
Part VIII

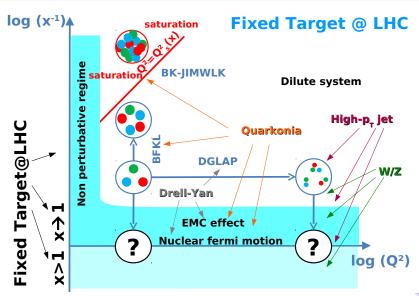
Backup slides











Gas target

C. Barschel, P. Lenisa, A. Nass, and E. Steffens, Adv. Hi. En. Phys. (2015) 463141

TABLE 1: Comparison of gas targets in storage rings with a hypothetical target for the proposed AFTER@LHC initiative [1, 2]. The target gas ¹H, ²D, or ³He is assumed to be spin polarized.

Storage ring	Particle	E _{max} [GeV]	Target type	L [m]	T [K]	L _{max} [1/cm ² s]	Remarks	Reference
HERA-e DESY (term. 2007)	e [±] pol.	27.6	Cell ¹ H, ² D, ³ He	0.4	100 25	$2.5 \cdot 10^{31} \\ 2.5 \cdot 10^{32}$	HERMES exp. 1995–2007	[9]
RHIC-p BNL	p pol.	250	Jet	_	_	$1.7 \cdot 10^{30}$	Absolute p polarimeter	[10]
COSY FZ Jülich	p, d pol.	3.77 T = 49.3 MeV	Cell ¹ H, ² D Cell ¹ H	0.4	300	$10^{29} \\ 2.75 \cdot 10^{29}$	ANKE exp. PAX exp.	[4, 5] [11]
LHC CERN (proposed)	p unpol. heavy ions	7,000 2,760 · A	Cell ${}^{1}H, {}^{2}D$ Xe $M \approx 131$	1.0	100 ≥100	$10^{33} \\ 10^{27} - 10^{28}$	Based on techn. of HERMES target	this paper

→ beam lifetime with $\mathcal{L}_{pp} = 10^{33} \text{cm}^{-2} \text{s}^{-1} = 10 \text{ nb}^{-1} \text{s}^{-1} \text{of } 2 \times 10^{6} \text{ s (or 23 days)}.$



→ Relevant parameters for existing and proposed polarized DY experiments.

S.J. Brodsky, F. Fleuret, C. Hadjidakis, JPL, Phys. Rep. 522 (2013) 239 V. Barone, F. Bradamante, A. Martin, Prog. Part. Nucl. Phys. 65 (2010) 267.

Experiment	particles	energy (GeV)	√s (GeV)	x_p^{\uparrow}	\mathcal{L} (nb ⁻¹ s ⁻¹)
AFTER	$p + p^{\uparrow}$	7000	115	0.01 ÷ 0.9	1
COMPASS	$\pi^{\pm} + p^{\uparrow}$	160	17.4	0.2 ÷ 0.3	2
COMPASS (low mass)	$\pi^{\pm} + p^{\uparrow}$	160	17.4	~ 0.05	2
P1039	$p + p^{\uparrow}$	120	15	0.1 ÷ 0.3	400-1000
P1027	$p^{\uparrow} + p$	120	15	$0.35 \div 0.85$	400-1000
RHIC	$p^{\uparrow} + p$	collider	500	$0.05 \div 0.1$	0.2
J-PARC	$p^{\uparrow} + p$	50	10	$0.5 \div 0.9$	1000
PANDA (low mass)	$\bar{p} + p^{\uparrow}$	15	5.5	$0.2 \div 0.4$	0.2
PAX	$p^{\uparrow} + \bar{p}$	collider	14	$0.1 \div 0.9$	0.002
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• AFTER could be the only project able to reach $x^{\uparrow} = 10^{-2}$ and $x^{\uparrow} > 0$, $4 = 10^{-2}$



• Gluon Sivers effect essentially unconstrained

D. Boer, C. Lorcé, C. Pisano, J. Zhou. Adv. Hi. En. Phys. (2015) 371396



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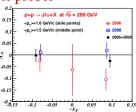
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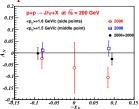


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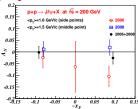
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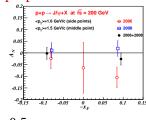
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y, y-jet, y - y
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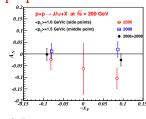
Gluon Sivers effect essentially unconstrained

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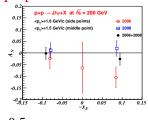
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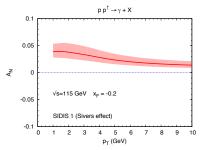
G. Mallot, Pacific Spin Symposium, 2015

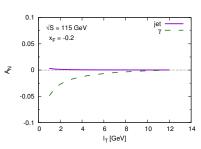
Further studies of the Sivers effect



• A_N^{γ} is predicted to have an opposite sign between the Generalised Parton Model (GPM) and the Collinear-Twist 3 (CT3) approach

GPM: M. Anselmino, U. D'Alesio, S. Melis. Adv.Hi.En.Phys. (2015) 475040 CT3: K. Kanazawa, Y. Koike, A. Metz, and D. Pitonyak. Adv.Hi.En.Phys. (2015) 257934.



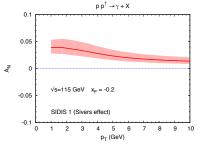


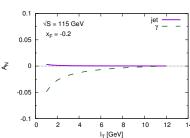
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- A_N^{π} : sign mismatch issue with $f_{1T}^{\perp,q}(x,\vec{k}_{\perp}^2)$ extracted from SIDIS
 - A_N^{jet} : complementary since no "contamination" (fragmentation Collins effect)
 - A_N^{π} should be measured at larger p_T

Instantaneous Luminosity:

$$\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A)/A$$

$$\Phi_{beam} = 2 \times 10^5 \text{ Pb s}^{-1}, \quad \ell = 1 \text{ cm (target thickness)}$$

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Target	ρ (g.cm ⁻³)	A	$\mathcal{L} \text{ (mb-1.s-1)} = \int \mathcal{L} \text{ (nb-1.yr-1)}$
1m Liq. H ₂	0.07	1	800
1m Liq. D ₂	0.16	2	1000
1cm Be	1.85	9	25
1cm Cu	8.96	64	17
1cm W	19.1	185	13
1cm Pb	11.35	207	7

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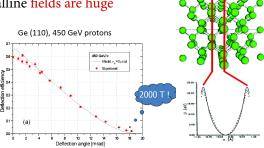
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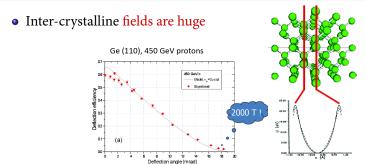
- Planned lumi for PHENIX Run15AuAu 2.8 nb⁻¹ (0.13 nb⁻¹ at 62 GeV)
- Nominal LHC lumi for PbPb 0.5 nb⁻¹

The beam extraction with a bent crystal

• Inter-crystalline fields are huge

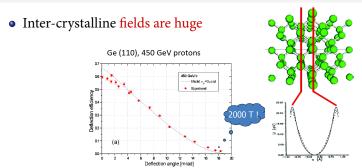


The beam extraction with a bent crystal



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Ge (110), 450 GeV protons

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Order of the state of the sta

- The channeling efficiency is high for a deflection of a few mrad
- One can extract a significant part of the beam loss $(10^9 p^+ s^{-1})$
- Simple and robust way to extract the most energetic beam ever:



The beam extraction: news

[S. Montesano, Physics at AFTER using LHC beams, ECT* Trento, Feb. 2013]

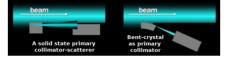
Goal : assess the possibility to use bent crystals as primary collimators in hadronic accelerators and colliders



UA9 installation in the SPS

Prototype crystal collimation system at SPS:

- local beam loss reduction (5÷20x reduction for proton beam)
- beam loss map show average loss reduction in the entire SPS ring
- halo extraction efficiency 70÷80% for protons (50÷70% for Pb)

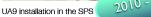


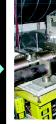
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Towards an installation in the LHC : propose and install during LSI a min. number of devices

• 2 crystals

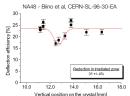
Long term plan is ambitious : propose a collimation system based on bent crystals for the upgrade of the current LHC collimation system

Simone Montesano - February 11th, 2013 - Physics at AFTER using the LHC beams

Crystal resistance to irradiation

- IHEP U-70 (Biryukov et al, NIMB 234, 23-30):
 - 70 GeV protons, 50 ms spills of 10¹⁴ protons every 9.6 s, several minutes irradiation
 - · equivalent to 2 nominal LHC bunches for 500 turns every 10 s
 - 5 mm silicon crystal, channeling efficiency unchanged
- * SPS North Area NA48 (Biino et al, CERN-SL-96-30-EA):
 - 450 GeV protons, 2.4 s spill of 5 x 10¹² protons every 14.4 s, one year irradiation, 2.4 x 10²⁰ protons/cm² in total,
 - · equivalent to several year of operation for a primary collimator in LHC
 - 10 x 50 x 0.9 mm³ silicon crystal, 0.8 x 0.3 mm² area irradiated, channeling efficiency reduced by 30%.
- HRMT16-UA9CRY (HiRadMat facility, November 2012);
 - * 440 GeV protons, up to 288 bunches in 7.2 μs , 1.1 x 10 11 protons per bunch (3 x 10 13 protons in total)
 - · energy deposition comparable to an asynchronous beam dump in LHC
 - 3 mm long silicon crystal, no damage to the crystal after accurate visual inspection, more tests planned to assess possible crystal lattice damage
 - accurate FLUKA simulation of energy deposition and residual dose







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- Extracted intensity: $5 \times 10^8 p^+ s^{-1}$ (1/2 the beam loss)

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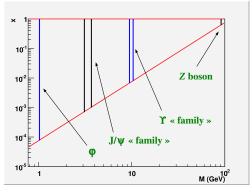
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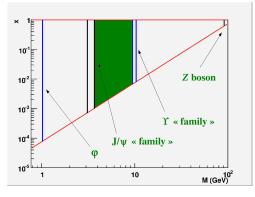
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• similar figures for the Pb-beam extraction

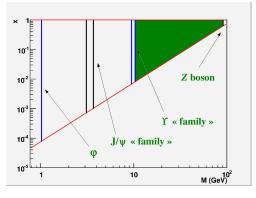
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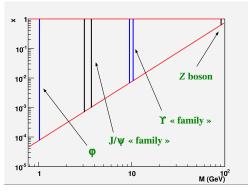


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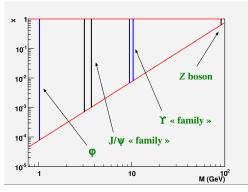
Note: $x_{target}(\equiv x_2) > x_{projectile}(\equiv x_1)$ "backward" region



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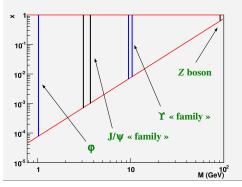
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To do: to look at the rates to see how competitive this will be

Our idea is not completely new

Nuclear Instruments and Methods in Physics Research A 333 (1993) 125-135 North-Holland NUCLEAR
INSTRUMENTS
& METHODS
IN PHYSICS
RESEARCH
SectionA

LHB, a fixed target experiment at LHC to measure CP violation in B mesons

Flavio Costantini

University of Pisa and INFN, Italy

A fixed target experiment at LHC to measure CP violation in B mesons is presented. A description of the proposed apparatus is given together with its sensitivity on the CP violation asymmetry measurement for the two benchmark decay channels $B^0 \rightarrow J/\psi + K_s^0$, $B^0 \rightarrow \pi^+\pi^-$. The possibility of obtaining an extracted LHC beam hinges on channeling in a bent silicon crystal. Recent results on beam extraction efficiencies measured at CERN SPS based on this technique are presented.

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1. Introduction

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This paper presents a fixed target experiment to measure CP violation in the B system based on the possibility of extracting the 8 TeV LHC proton beam using a bent silicon crystal [4]. A 10% extraction efficiency of the LHC beam halo will give an extracted beam intensity of about 10^8 protons/s allowing the production of as many as 10^{10} BB pairs per year, i.e. about two orders of magnitude more than what could be produced by an e⁺e⁻ asymmetric B factory with 10^{34} cm⁻²s⁻¹ luminosity [5].



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- After a year, one simply moves the crystal by less than one mm ...

Accessing the large *x* glue with quarkonia:

PYTHIA simulation $\sigma(y)$ / $\sigma(y=0.4)$ statistics for one month 5% acceptance considered

Statistical relative uncertainty Large statistics allow to access very backward region

Gluon uncertainty from MSTWPDF

- only for the gluon content of the target
- assuming

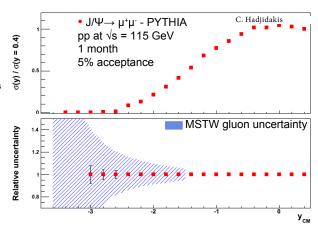
$$x_g = M_{J/\Psi}/\sqrt{s} e^{-yCM}$$

J/Ψ

$$y_{CM} \sim 0 \rightarrow x_g = 0.03$$

 $y_{CM} \sim -3.6 \rightarrow x_g = 1$

Y: larger x_g for same ycm $y_{CM} \sim 0 \rightarrow x_g = 0.08$ $y_{CM} \sim -2.4 \rightarrow x_g = 1$



⇒ Backward measurements allow to access large x gluon pdf

Assuming that we understand the quarkonium-production mechanisms

